# Experimental and theoretical investigations on the structure-properties interrelationship of the gadolinium-vanadate-germanate glasses 

Simona Rada • Ramona Chelcea • Eugen Culea

Received: 18 December 2009 / Accepted: 11 March 2010 /Published online: 19 April 2010
(C) Springer-Verlag 2010


#### Abstract

Glasses in the system $\mathrm{xGd}_{2} \mathrm{O}_{3} \cdot(100-\mathrm{x})\left[\mathrm{GeO}_{2} \cdot \mathrm{~V}_{2} \mathrm{O}_{5}\right]$ with $0 \leq x \leq 20 \mathrm{~mol} \%$ have been prepared from the melt quenching method. In this paper, we investigated changes in germanium coordination number in gadolinium-vanadategermanate glasses through molar volume analysis, measurements of densities, investigations of FTIR and UV-VIS spectroscopy, calculations of density functional theory (DFT). Analyzing the structural changes resulted from the IR spectra we found that the gadolinium ions have a pronounced affinity toward $\left[\mathrm{VO}_{4}\right]$ structural units which contain non-bridging oxygens necessary for the charge compensation. The introduction of the excess of oxygen yields the formation of $\left[\mathrm{VO}_{5}\right]$ structural units. This attains maximum value at $5 \mathrm{~mol} \% \mathrm{Gd}_{2} \mathrm{O}_{3}$, in agreement with the density measurements. Further, the addition of the surplus of oxygen implies the transformation of $\left[\mathrm{VO}_{5}\right]$ to $\left[\mathrm{VO}_{4}\right]$ structural units and the formation of $\mathrm{VO}_{4}^{-3}$ orthovanadate structural units. The UV-VIS spectra show a broad UV absorption band located in the $300-500 \mathrm{~nm}$ region. These bands are assumed to originate from the combination of vanadium ions possibly present in the three states of valence. The presence of $\mathrm{Ge}-\mathrm{Ge}$ wrong bonds attains its maximum values in the samples with $\mathrm{x}=5$ and $15 \mathrm{~mol} \%$ $\mathrm{Gd}_{2} \mathrm{O}_{3}$ (bands centered in the $250-300 \mathrm{~nm}$ range). DFT calculations show the massive vibrations of the $\left[\mathrm{VO}_{\mathrm{n}}\right]$ structural units coupled with each other via $\left[\mathrm{GeO}_{6}\right]$ and $\left[\mathrm{GeO}_{4}\right]$ structural units. This leads to the splitting of the


[^0]bridge modes and a multiplication of the number of these bands.

Keywords DFT calculations • FT-IR spectroscopy• Gadolinium-vanadate-germanate glasses. UV-VIS spectroscopy

## Introduction

Germanate glasses doped with rare-earth ions have been investigated extensively in the past because their physicochemical, optical and spectroscopic properties are advantageous for optoelectronic applications [1-4].

The structural chemistry of crystalline oxides such as $\mathrm{V}_{2} \mathrm{O}_{5}$ revealed that the structure of $\mathrm{V}_{2} \mathrm{O}_{5}$ is built of octahedrons, where the pentavalent vanadium has five-coordination with oxygen atoms. Although $\mathrm{V}^{+5}$ thus behaves like $\mathrm{P}^{+5}$ in forming the ions $\mathrm{VO}_{4}^{-3}$ and $\mathrm{V}_{2} \mathrm{O}_{7}^{-4}$ it shows a much greater tendency than P to form condensed oxy-ions.

On the other hand, the $\mathrm{V}_{2} \mathrm{O}_{5}$-rich glasses in which $\mathrm{V}_{2} \mathrm{O}_{5}$ acts as the network former have the network structure mainly consisting of corner-sharing branched $\left[\mathrm{VO}_{4}\right]$ tetrahedral of the same structural units as found in phosphate glasses. The network structure was reported to be made up of unaffected [ $\mathrm{VO}_{5}$ ] units as in vitreous $\mathrm{V}_{2} \mathrm{O}_{5}$ and affected [VO5] units with alkaline earth ions in contrast to the vanadate glasses formed by conventional network formers in which only unaffected $\left[\mathrm{VO}_{5}\right]$ units are present [4-9]. These glasses are known to contain $\mathrm{V}^{+4}$ and $\mathrm{V}^{+5}$ ions where the electrical conduction was attributed to the hopping of $3 \mathrm{~d}^{1}$ unpaired electron from $\mathrm{V}^{+4}$ to $\mathrm{V}^{+5}$ site which induces a polarization of the vanadium ion around it and form a polaron [10-14].

## Experimental and computational procedure

The gadolinium-vanadate-germanate glasses with composition $\mathrm{xGd}_{2} \mathrm{O}_{3} \cdot(100-\mathrm{x})\left[\mathrm{GeO}_{2} \cdot \mathrm{~V}_{2} \mathrm{O}_{5}\right]$ where $\mathrm{x}=0,5,10,15$, $20 \%$ moli were prepared by taking appropriate amounts of reagents grade $\mathrm{GeO}_{2}, \mathrm{~V}_{2} \mathrm{O}_{5}$ and $\mathrm{Gd}_{2} \mathrm{O}_{3}$. After melting at $1000^{\circ} \mathrm{C}$ for about 10 min in a corundum crucibles and introduce in an electric furnace, the melt was quenched rapidly to room temperature.

The samples were analyzed by means of X-ray diffraction using a XRD-6000 Shimadzu diffractometer, with a monochromator of graphite for the $\mathrm{Cu}-\mathrm{K} \alpha$ radiation $(\lambda=$ $1.54 \AA$ ) at room temperature.

The structure of the glasses was investigated by infrared spectroscopy using the KBr pellet technique. The IR spectra were recorded in the range $400-1200 \mathrm{~cm}^{-1}$ using a JASCO FTIR 6200 spectrophotometer.

The densities of the glass samples were measured accurately to the third decimal $\left( \pm 0,03 \mathrm{~g} \mathrm{~cm}^{-3}\right)$ by the displacement method using $\mathrm{CCl}_{4}$ as an immersion liquid.

The starting structures have been built using the graphical interface of Spartan'04 [15] and preoptimized by molecular mechanics. Optimizations were continued at DFT level (B3PW91/CEP-4G/ECP) using the Gaussian'03 package of programs [16].

It should be noticed that only the broken bonds at the model boundary were terminated by hydrogen atoms. The positions of boundary atoms were frozen during a calculation and the coordinates of internal atoms were optimized, to model the active fragment flexibility and its incorporation into the bulk.

## Results and discussion

Density, molar volume and oxygen packing density

Density is a powerful tool capable of exploring the changes in the structure of glasses. Density is affected by the structural softening/compactness, change in geometrical configuration, coordination number, cross-link density and dimension of interstitial spaces of the glasses.

Figure 1 shows the compositional evolution of the density of gadolinium-vanadate-germanate glasses. This figure shows the presence of density maxima at $\mathrm{x}=5$ and $20 \mathrm{~mol} \% \mathrm{Gd}_{2} \mathrm{O}_{3}$. The changes in molar volume were particularly pronounced for $\mathrm{x}=0 \%$ and appeared to be minimal at $\mathrm{x}=5$ and $20 \mathrm{~mol} \%$. We conclude that with increasing of the $\mathrm{Gd}_{2} \mathrm{O}_{3}$ content up to $5 \mathrm{~mol} \%$, the local vanadium environment in the glass changes from tetrahedral to pentagonal geometry due to the fact that the coordination number of vanadium atoms changes from 4 to 5 . Further the addition of the modifier gadolinium oxide


Fig. 1 Gadolinium oxide composition dependence on (a) density, (b) molar volume, $\mathrm{V}_{\mathrm{m}}$ and (c) the oxygen packing density, $\mathrm{d}_{\mathrm{O}}$, for $\mathrm{xGd}_{2} \mathrm{O}_{3} \cdot(100-\mathrm{x})\left[\mathrm{GeO}_{2} \cdot \mathrm{~V}_{2} \mathrm{O}_{5}\right]$ samples
to the vanadium germanate glass network introduces a surplus oxygen into the vitreous network and the need to accommodate the glasses network. The additional oxygen may be incorporated into the network as non-bridging oxygen by breaking oxygen bridges and in these conditions the four-coordinated germanium atoms attain maximum value.

The oxygen packing density is the ratio of measured density per molecular weight and number of oxygen ions in a formula unit [17]. For $x \leq 5 \mathrm{~mol} \%$ and $\mathrm{x}=20 \%$ moli the oxygen packing density increases which clearly indicates that the glassy network becomes tightly packed when more $\mathrm{Gd}_{2} \mathrm{O}_{3}$ is introduced to the vitreous matrix.

In the case of pure $\mathrm{V}_{2} \mathrm{O}_{5}$ glass it has been reported [14] that $\mathrm{V}^{+5}$ ions exhibit both four and fivefold coordination states, depending on the sample preparation conditions. The IR spectrum of pure crystalline and amorphous $\mathrm{V}_{2} \mathrm{O}_{5}$ is characterized by the intense band in the $980-1020 \mathrm{~cm}^{-1}$ region, related to vibration of isolated $\mathrm{V}=\mathrm{O}$ vanadyl groups [18, 19] while the band located at about $950-$ $970 \mathrm{~cm}^{-1}$ was attributed to $\left[\mathrm{VO}_{4}\right]$ units [20-23]. If $\mathrm{V}=\mathrm{O}$ mode position at around $1020 \mathrm{~cm}^{-1}$ is preserved it can be concluded that $\mathrm{V}=\mathrm{O}$ bond is not directly influenced and the coordination number and symmetry of the $\left[\mathrm{VO}_{4}\right]$ and [ $\mathrm{VO}_{5}$ ] structural units do not change.

## FTIR spectroscopy

FTIR spectra obtained for the samples are shown in Fig. 2. The examinations of these spectra show that the increase of $\mathrm{Gd}_{2} \mathrm{O}_{3}$ content modify the characteristic IR bands as follows:
i) It can be seen that the IR band in the wavenumber region of $\sim 800 \mathrm{~cm}^{-1}$ is found for all samples and its


Fig. 2 FTIR spectra of the $\mathrm{xGd}_{2} \mathrm{O}_{3} \cdot(100-\mathrm{x})\left[\mathrm{GeO}_{2} \cdot \mathrm{~V}_{2} \mathrm{O}_{5}\right]$ samples with $\mathrm{x}=0-20 \% \mathrm{Gd}_{2} \mathrm{O}_{3}$
intensity attains maximum value at $20 \mathrm{moli} \% \mathrm{Gd}_{2} \mathrm{O}_{3}$. This band can be due to the $\mathrm{Ge}-\mathrm{O}$ stretching vibrations of the $\left[\mathrm{GeO}_{4}\right]$ tetrahedral structural units involving nonbridging oxygen [10, 11]. This suggests that the glass network modification has taken place mainly in the vanadate part whereas the germanate part remained unmodified and its network consists mainly from the $\left[\mathrm{GeO}_{4}\right]$ structural units with interconnected through Ge-O-Ge bridges in $\left[\mathrm{GeO}_{4}\right]$ structural units.
ii) The bands located at about $480 \mathrm{~cm}^{-1}$ and $530 \mathrm{~cm}^{-1}$ attain their maximum values at $\mathrm{x}=10 \mathrm{~mol} \% \mathrm{Gd}_{2} \mathrm{O}_{3}$. These bands are attributed to the angular deformation vibration of the V-O bond [24]. This higher capacity of migration of the gadolinium ions inside the glass can be explained by a process of conversion of the $\left[\mathrm{VO}_{5}\right]$ to $\left[\mathrm{VO}_{4}\right]$ structural units.
iii) With the increase of $\mathrm{Gd}_{2} \mathrm{O}_{3}$ content, a sharp decreasing trend was observed in wavenumber of the bands centered at about $1025 \mathrm{~cm}^{-1}$, which belong to vibration of isolated $\mathrm{V}=\mathrm{O}$ bond. This may be attributed to the electrostatic field of the strongly polarizing $\mathrm{Gd}^{+3}$ ions. The increase of $\mathrm{Gd}^{+3}$ content leads to strengthening the electron cloud around oxygen in $\left[\mathrm{VO}_{\mathrm{n}}\right]$ units, and consequently causes a shift of stretching vibration of V-O-Ge and/or V-O-Gd to lower wavenumber. The intensity of this band attains its maximum value at $5 \% \mathrm{Gd}_{2} \mathrm{O}_{3}$.

The shifting of the high-frequency band of the $\mathrm{V}=\mathrm{O}$ bond toward lower wavenumber can be attributed to changes in the structure of vanadate network produced by the addition of $\mathrm{Gd}_{2} \mathrm{O}_{3}$. When $\mathrm{Gd}_{2} \mathrm{O}_{3}$ is introduced, $\mathrm{Gd}^{+3}$ ions go either to interstitial or substitutional sites in vanadate chains. This would influence the isolated $\mathrm{V}=\mathrm{O}$ bonds and thus affect the
vibrations of $\mathrm{V}=\mathrm{O}$ bonds depending upon the position of the gadolinium ions. If gadolinium ions occupy substitutional sites, the V-O-V bonds will break up and form V-O-Gd bridges which restrict the influence of gadolinium ions on $\mathrm{V}=$ O bonds, yielding to preserve the position of the vibrational frequency at $1025 \mathrm{~cm}^{-1}$ in the IR spectrum.
iv) The bands located at about $980 \mathrm{~cm}^{-1}$ is attributable not only to the $\left[\mathrm{VO}_{5}\right]$ structural units but also to the branched $\left[\mathrm{VO}_{4}\right]$ structural units having one $\mathrm{V}=\mathrm{O}$ double bond. In accordance with the mechanism suggested above the gadolinium ions rest between vanadates chains and their influence on the $\mathrm{V}=\mathrm{O}$ bond is important. This band shows a maximum intensity at $\mathrm{x}=5 \mathrm{~mol} \%$.
v) v) An increasing trend was observed in the strength of the bands centered at $\sim 870 \mathrm{~cm}^{-1}$. This band is usually assigned to the stretching vibrations of $\mathrm{VO}_{4}{ }^{-3}$ orthovanadate structural units.

The added $\mathrm{Gd}_{2} \mathrm{O}_{3}$ also gives rise to the formation of nonbridging oxygens, thereby creating $\mathrm{VO}_{4}{ }^{-3}$ units. It is obvious that the oxygen, which becomes non-bridging and acquires a negative charge, will move closer to the connected vanadium, consequently reducing the positive charge on the vanadium and thereby resulting in a decrease in the binding of other oxygens attached to this particular ion; the length of the $\mathrm{V}=\mathrm{O}$ bond therefore increases. Also, due to the irregular and random distribution of atoms in the glass structure, the gadolinium ions take positions interstitially that are more symmetrical among other units.

Analyzing the structural changes resulted from the IR spectra, we found that the gadolinium ions have a pronounced affinity toward $\left[\mathrm{VO}_{4}\right]$ structural units which contain non-bridging oxygens necessary for the charge compensation. The introduction of the excess of oxygen yields the formation of $\left[\mathrm{VO}_{5}\right]$ structural units. This attains maximum value at $5 \mathrm{~mol} \% \mathrm{Gd}_{2} \mathrm{O}_{3}$, in agreement with the density measurements. Further, the addition of the surplus of oxygen implies the transformation of $\left[\mathrm{VO}_{5}\right]$ to $\left[\mathrm{VO}_{4}\right]$ structural units and the formation of $\mathrm{VO}_{4}^{-3}$ orthovanadate structural units. The formation of free orthovanadate units creates smaller network cavities, the increase of local densification and density (for sample with $\mathrm{x}=20 \mathrm{moli} \%$ ).

## UV-VIS spectroscopy

The UV-VIS absorption spectra of $\mathrm{xGd}_{2} \mathrm{O}_{3} \cdot(100-\mathrm{x})$ $\left[\mathrm{GeO}_{2} \cdot \mathrm{~V}_{2} \mathrm{O}_{5}\right]$ glasses are shown in Fig. 3. The examinations of these spectra show that the characteristic UV-VIS bands are modified namely:
i) For samples with $\mathrm{x} \geq 5 \mathrm{~mol} \% \mathrm{Gd}_{2} \mathrm{O}_{3}$, it follows from the spectrum that the glass matrix absorption begins at


Fig. 3 UV-VIS absorption spectra of samples in function of gadolinium oxide content
about $250 \mathrm{~cm}^{-1}$. These bands are very sharp in the lead germanate with $\mathrm{x}=5$ and $15 \mathrm{~mol} \%$, and shifted with increasing of gadolinium oxide content to longer wavenumbers ( 275 nm ). All these bands can be assigned to the Ge-Ge wrong bonds [25].

This shifting can be interpreted considering the chemical bonding property of the ligand (oxygen) that exerts influences on the positions the absorption peaks due to the so-called nephelauxetic effect (meaning cloud expanding effect) as defined by Jorgensen [26]. As a consequence of this effect when a paramagnetic ion (gadolinium ion) is introduced into glass, parameters such as the electrostatic interaction of the electron, spin-orbit interaction and potential at the ion site due to its environment the Hamiltonian expression are reduced from their free-ion
values [27]. The chemical bonding property of the ligand or the nephelauxetic effect can be elucided in terms of electronegativity because it reflects the degree of covalent bonding characteristics of different hosts of $\mathrm{Gd}^{+3}$ ions [28].

The changes in the features of the absorption bands can be explained as a consequence of the appearance of additional absorption shoulder due to photoinduced color centers in the glass such as the formation of germanategadolinium paramagnetic defect centers in the glasses.

The data above seem to confirm that the presence of gadolinium ions, well known as glass network modifiers and necessary for the ion-exchange, modify the structure of the glasses with $\mathrm{x}=5$ and $15 \mathrm{~mol} \% \mathrm{Gd}_{2} \mathrm{O}_{3}$. A striking evidence is that the bands situated in the $250-300 \mathrm{~nm}$ region, is not detectable if gadolinium ions are present in the amounts corresponding to $\mathrm{x}=20 \mathrm{~mol} \%$. This could be due to the formation of a gadolinium germanate-like structure, which is able to eliminate or hide the presence of $\mathrm{Ge}-\mathrm{Ge}$ wrong bonds absorbing at $250 \mathrm{~cm}^{-1}$. If any $\mathrm{Ge}-\mathrm{Ge}$ wrong bonds are present in these glasses, they do not absorb at $250 \mathrm{~cm}^{-1}$ anymore.
ii) It follows from the spectrum that all the glasses samples show a strong absorption extending from about 300 to 530 nm . This broad UV absorption is assumed to originate from the combination of vanadium ions possibly present in the three valences. Trivalent vanadium ions exhibit a characteristic UV band at $350-400 \mathrm{~nm}$, vanadyl tetravalent ions possess an absorption at about 420 nm and pentavalent ions are known to belong to $\mathrm{d}^{0}$ configuration and exhibit only UV bands.
The bands located at about 360 nm correspond to the $\mathrm{V}^{+5}$ ions. These results show that the incorporation of $\mathrm{Gd}^{+3}$ ions in the host matrix have directly influenced $\mathrm{V}=\mathrm{O}$ bonds. The $\mathrm{V}=\mathrm{O}$ bonds are not conserved. The UV-VIS

Fig. 4 The optimized structure of the $\mathrm{GeO}_{2} \cdot \mathrm{~V}_{2} \mathrm{O}_{5}$ glasses was used to perform the DFT computations



Fig. 5 The experimental and theoretical IR spectra for the $\mathrm{GeO}_{2} \cdot \mathrm{~V}_{2} \mathrm{O}_{5}$ matrix network
spectrum is dominated by the charge transfer of the type $\mathrm{O}^{-2} \rightarrow \mathrm{~V}^{+5}[29,30]$.
iii) For the sample with $x=1 \%$ moli $\mathrm{Gd}_{2} \mathrm{O}_{3}$, our results reveal a low resolution of the bands located in the range from $600-800 \mathrm{~nm}$. These are due to the $\mathrm{V}^{+4}$ ions.

The presence of $\mathrm{Ge}-\mathrm{Ge}$ wrong bonds attains its maximum values in the samples with $\mathrm{x}=5$ and $15 \mathrm{~mol} \% \mathrm{Gd}_{2} \mathrm{O}_{3}$. The UV-VIS spectra of all samples are dominated by the charge transfer of the type $\mathrm{O}^{-2} \rightarrow \mathrm{~V}^{+5}$ (the $\mathrm{V}^{+5}$ ions have $\mathrm{d}^{0}$ configuration) at about $\sim 380 \mathrm{~nm}$ (with an edge energy of about 3.18 eV ) [29, 30]. By increasing the gadolinium
oxide content up $20 \mathrm{~mol} \%, \mathrm{Gd}^{+3}$ ions rest between vanadates chains may affect the isolated $\mathrm{V}=\mathrm{O}$ bonds.

## DFT calculations

These data will be used in the present research in order to compute a possible structural model of the $\mathrm{GeO}_{2} \cdot \mathrm{~V}_{2} \mathrm{O}_{5}$ glasses network. Similar methodology has previously been reported to study other glasses [31-34]. The study on structural modifications of the vitreous network and the equilibrium geometry were found by optimization (Fig. 4).

There are some sites in our model (Fig. 4), namely:
i) Germanium atom is in hexagonal coordination. The Ge-O bonds are subdivided into four groups: a bond has shorter interatomic distance ( $1.78 \AA$ ), two bonds have short interatomic distances $(1.82 \AA)$, one has intermediate interatomic distances $(1.88 \AA)$ and two have longer interatomic distances ( $2.65 \AA$ ). Considering all six bonds, the $\left[\mathrm{GeO}_{6}\right]$ polyhedron is irregular with four $\mathrm{Ge}-\mathrm{O}$ bond lengths ranging from 1.78 to $1.88 \AA$ and with two oxygen atoms remaining relatively far away from germanium cations ( $2.65 \AA$ ). These bonds are too large to be accepted as chemical bonds in the usual sense.
Therefore, if only oxygen atoms with Ge-O distances less than $1.88 \AA$ are taken into account, the coordination polyhedron of the germanium is equal to four. However, from crystallographic point of view, coordination polyhedron of the germanium atoms is equal to six.

This shows that there is instability between the nonequivalent $\mathrm{Ge}-\mathrm{O}$ bonds in the octahedron. In essence, there is a large

Fig. 6 The distribution of the electronic states of the HOMO $\left(\mathrm{E}_{\mathrm{HOMO}}=-9.90 \mathrm{eV}\right)$, HOMO-1, LUMO ( $\mathrm{E}_{\text {LUMO }}=-2.59 \mathrm{eV}$ ) and LUMO+1 of the proposed model for $\mathrm{GeO}_{2} \cdot \mathrm{~V}_{2} \mathrm{O}_{5}$ glasses network

displacement of the germanium atom from the center of the polyhedron, whose strongly asymmetry depends on the manner of connection with the surrounding polyhedron. This deformation will be expressed more clearly when one or two of the oxygen atoms in the vertex are not bridged, which is typical in the formation of the vitreous matrix.
ii) The germanium site is coordinated to four O atoms. The Ge-O bonds are subdivided into two groups: two bonds have short interatomic distances (1.77-1.79 $\AA$ ) and two have longer interatomic distances (1.80$1.83 \AA$ ).
iii) The vanadium ions are distributed into two crystallographic sites: the $\left[\mathrm{VO}_{4}\right]$ tetrahedrons and $\left[\mathrm{VO}_{5}\right]$ square pyramids. The $\left[\mathrm{VO}_{4}\right]$ tetrahedral units are higher distorted around the vanadium center (the vanadiumoxygen distances are ranging from 1.62 to $2.02 \AA$ ). The $\left[\mathrm{VO}_{5}\right]$ square pyramids are considerably distorted around vanadium center and the V-O bond distances are ranging from 1.64 to $3.78 \AA$. Such behavior was reported for two-dimensional layered vanadate compounds [35, 36].
The vibrational properties of these glasses inferred from quantum chemical simulation are interpreted through a comparison with the experimental spectrum (Fig. 5). The evolution of vibrational spectrum of the proposed model is important for understanding the broadening-effect of the glasses from the experimental FTIR spectrum. Our results show that the Ge-O and V-O stretching vibration region of the proposed model is similar to the same region of the glass.

The massive vibrations of the $\left[\mathrm{VO}_{\mathrm{n}}\right]$ structural units can be coupled with each other via $\left[\mathrm{GeO}_{6}\right]$ and $\left[\mathrm{GeO}_{4}\right]$ structural units. This leads to the splitting of the bridge modes and a multiplication of the number of these bands.

The distribution of the electronic states of the HOMO, HOMO-1, LUMO and LUMO+1 can be seen in Fig. 6. An interesting finding in these systems is that:
i) The HOMO and HOMO-1 give character of electron donor for the $\left[\mathrm{VO}_{4}\right],\left[\mathrm{VO}_{5}\right]$ structural units of the vanadate glasses network and the $\left[\mathrm{GeO}_{6}\right]$ structural units. The HOMO energy, $\mathrm{E}_{\text {номо }}$ is -9.90 eV .
ii) The LUMO and LUMO+1 give the character of electron acceptor for the $\left[\mathrm{GeO}_{4}\right]$ structural units of the germanate network and the $\left[\mathrm{VO}_{\mathrm{n}}\right]$ structural units of the vanadate network. The LUMO energy, $\mathrm{E}_{\mathrm{LUMO}}$ is -2.59 eV .

There is a change transfer between the vanadium atoms tetra- and penta-coordinated. Also, a charge transfer was yielded between germanium atoms and vanadate network. This can be explained considering that the vanadate network is flexible to form the appropriate coordination environments with structural units of opposite charge such as gadolinium ions.

## Conclusions

Homogeneous glasses of $\mathrm{xGd}_{2} \mathrm{O}_{3} \cdot(100-\mathrm{x})\left[\mathrm{GeO}_{2} \cdot \mathrm{~V}_{2} \mathrm{O}_{5}\right]$ system were obtained within $0 \leq x \leq 20 \mathrm{~mol} \%$. In this paper, we investigated changes in structural properties of the gadolinium-vanadate-germanate glasses through molar volume analysis, measurements of densities, investigations of FTIR and UV-VIS spectroscopy and calculations of density functional theory (DFT).

Our results show that the gadolinium ions have a pronounced affinity toward [ $\mathrm{VO}_{4}$ ] structural units which contain non-bridging oxygens necessary for the charge compensation. The formation of $\left[\mathrm{VO}_{5}\right]$ structural units attains maximum value at $5 \mathrm{~mol} \% \mathrm{Gd}_{2} \mathrm{O}_{3}$, in agreement with the density measurements. Further, the addition of the surplus of oxygen implies the transformation of $\left[\mathrm{VO}_{5}\right]$ to $\left[\mathrm{VO}_{4}\right]$ structural units and the formation of $\mathrm{VO}_{4}^{-3}$ orthovanadate structural units.

The broader UV absorption bands located in the 300500 nm region corresponds to the vanadium ions present in the three states of valence. In the sample with $x=5$ and $15 \mathrm{~mol} \% \mathrm{Gd}_{2} \mathrm{O}_{3}$, the $\mathrm{Ge}-\mathrm{Ge}$ wrong bonds attains its maximum values. DFT calculations show the massive vibrations of the $\left[\mathrm{VO}_{n}\right]$ structural units coupled with each other via $\left[\mathrm{GeO}_{6}\right]$ and $\left[\mathrm{GeO}_{4}\right]$ structural units.

Acknowledgments The financial support of the Ministry of Education and Research of Romania-National University Research Council (CNCSIS, PN II-IDEI 183/2008, contract number 476/2009) is gratefully acknowledged by the authors.

## References

1. Guinhos FC, Nobrega PC, Santa-Cruz PA (2001) J Alloys Compd 323\&324:358-361
2. Bueno LA, Melnikov P, Messaddeg Y, Ribeiro SJL (1999) J NonCryst Solids 247:87-91
3. Klimesz B, Dominiak-Dzik G, Lisiecki R, Ryba-Romanowski W (2008) J Non-Cryst Solids 354:515-520
4. Wright AC, Yarker CA, Johnson PAV (1985) J Non-Cryst Solids 76:333-350
5. Sen S, Ghosh A (2001) J Phys Condens Matter 13:1979-1986
6. Sen S, Ghost A (1999) J Non-Cryst Solids 258:29-33
7. Dimitriev V, Dimitrov Y, Arnaudov M, Topalov D (1983) J NonCryst Solids 57:147-156
8. Sayer M, Mansingh A (1972) Phys Rev B 6:4629-4643
9. Mott NF (1968) J Non-Cryst Solids 1:1-18
10. Rada M, Culea E, Rada S, Maties V, Pascuta P (2010) J Mater Sci 45(6):1487-1494
11. Chung CH, Mackenzie JD (1980) J Non-Cryst Solids 42:357-370
12. Owen AE (1977) J Non-Cryst Solids 25:370-423
13. Sen S, Ghosh A (1999) Phys Rev B 60(22):15143-15149
14. Mendialdua J, Casanova R, Barbaux Y (1995) J Electron Spectrosc Relat Phenom 71:249-261
15. Spartan'04, Wavefunction Inc, 18401 Von Karman Avenue, Suite 370 Irvine, CA 92612
16. Frisch MJ, Trucks GW, Schlegel HB, Scuseria GE, Robb MA, Cheeseman JR, Montgomery JA, Vreven T, Kudin KN, Burant

JC, Millam JM, Iyengar SS, Tomasi J, Barone V, Mennucci B, Cossi M, Scalmani G, Rega N, Petersson GA, Nakatsuji H, Hada M, Ehara M, Toyota K, Fukuda R, Hasegawa J, Ishida M, Nakajima T, Honda Y, Kitao O, Nakai H, Klene M, Li X, Knox JE, Hratchian HP, Cross JB, Adamo C, Jaramillo J, Gomperts R, Stratmann RE, Yazyev O, Austin AJ, Cammi R, Pomelli C, Ochterski JW, Ayala PY, Morokuma K, Voth GA, Salvador P, Dannenberg JJ, Zakrzewski VG, Dapprich S, Daniels AD, Strain MC, Farkas O, Malick DK, Rabuck AD, Raghavachari K, Foresman JB, Ortiz JV, Cui Q, Baboul AG, Clifford S, Cioslowski J, Stefanov BB, Liu G, Liashenko A, Piskorz P, Komaromi I, Martin RL, Fox DJ, Keith T, Al-Laham MA, Peng CY, Nanayakkara A, Challacombe M, Gill PMW, Johnson B, Chen W, Wong MW, Gonzalez C, Pople JA (2003) Gaussian 03, revision A1. Gaussian Inc, Pittsburgh
17. Miyata H, Fujii K, Ono T, Kubokawa Y, Ohno T, Hatayama F (1987) J Chem Soc Faraday Trans 83:675-685
18. Rada S, Bosca M, Culea E, Rada M, Dan V, Maties V (2009) Struct Chem 20:801-805
19. Dimitrov V (1987) J Solid State Chem 66:256
20. Khattak GD, Tabet N, Wenger LE (2005) Phys Rev B 72:104203104215
21. Rada S, Culea M, Rada M, Culea E (2008) J Mater Sci 43 (18):6122-6125
22. Rada S, Culea E, Rus V, Pica M, Culea M (2008) J Mater Sci 43:3713-3716
23. Manara D, Grandjean A, Pinet O, Dussossoy JL, Neuville DR (2007) J Non-Cryst Solids 353:12-23
24. Mandal S, Hazra S, Das D, Ghosh A (1995) J Non-Cryst Solids 183:315-319
25. Atkins RM, Mizrahi V (1992) Electron Lett 28:1743-1744
26. Jorgensen CK, Pappalardo R, Schmidtker HH (1963) J Chem Phys 39:1422-1430
27. Patek K (1970) Glass lasers. Butterworth, London
28. Katzin L, Barnet ML (1964) J Phys Chem 68:3779-3785
29. Keller DE, Visser T, Soulimani F, Konigsberger DC, Weckhuysen BM (2007) Vibrat Spectrosc 43:140-151
30. Gao X, Wachs IE (2000) J Phys Chem B 104:1261-1268
31. Rada S, Culea M, Culea E (2008) J Phys Chem A 112(44):1125111255
32. Rada S, Culea M, Neumann M, Culea E (2008) Chem Phys Lett 460:196-199
33. Rada S, Culea E (2009) J Mol Struct 929:141-148
34. Rada S, Ristoiu T, Rada M, Coroiu I, Maties V, Culea E (2010) Mater Res Bull 45:69-73
35. Sun C, Wang E, Xiao D, An D, Xu L (2007) J Mol Struct 840: 53-58
36. Zhou Y, Qiao H (2007) Inorg Chem Commun 10:1318-1320


[^0]:    S. Rada $(\bowtie) \cdot$ R. Chelcea $\cdot$ E. Culea

    Department of Physics, Technical University of Cluj-Napoca, 400020 Cluj-Napoca, Romania
    e-mail: Simona.Rada@phys.utcluj.ro

